



# Stability and recursive solutions in Hamiltonian PDEs

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# Non linear PDE's

Non linear PDE's are ubiquitous in Mathematics and Physics. Some of the most famous examples as the Euler or Navier Stokes equations come from Hydrodynamics, as well as the Korteweg de Vries, Non Linear Schrödinger, Camassa Holm etc.

## Why Non-linearity?



Non linearity describes mathematically **self interactions** of waves, fields etc..

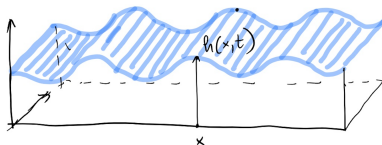
# Non linear PDE's

Here are some well known examples in  $1d$  :

$$(NLS) \quad u_t = -i(u_{xx} + |u|^2 u),$$

$$(KdV) \quad u_t = u_{xxx} + uu_x,$$

$$(CH) \quad u_t - u_{xxt} - u_x + u_{xxx} + uu_{xxx} + 3uu_x - 2u_x u_{xx} = 0$$



These are in fact all **integrable PDEs** which are widely studied as models in hydrodynamics.

The phenomena that I shall discuss do not rely on an integrable structure! (you could add a small perturbation!)

One is interested also in modeling waves in higher dimensional domains

$$\begin{cases} -i\partial_t u + \Delta u = |u|^{2p}u, \\ u = u(t, x), \quad x \in D \end{cases} \quad (NLS)$$

$D$  can be a domain in  $\mathbb{R}^n$ , or a Riemannian manifold..

I will concentrate on **PDEs defined on compact manifolds without boundary** , mainly **tori**

$$\mathbb{T}^n := \mathbb{R}^n / (2\pi\mathbb{Z})^n, \quad u(\cdot, x) \text{ is } 2\pi \text{ periodic in each } x_i$$

When treating non-linear phenomena one should expect a very complex behavior with instability, chaos etc. (as well known to meteorologists!).

Many interesting models have some symmetry such as:

**Hamiltonian or Reversible structure**

and this leads to a coexistence of **regular** and **chaotic** orbits

I will concentrate on regular behavior.

consider PDEs on tori  $\mathbb{T}^n$  close to an elliptic fixed point and study existence and stability of invariant tori taking a dynamical system point of view.

# Dynamical systems in $\infty$ dimension.

we consider a **non-linear dynamical system**:

$$\dot{u} = F(u), \quad u \in V$$

where

- $V$  is a vector space (in our case a scale of Hilbert spaces)
- $F$  is a non-linear functional from  $V$  in itself.

# PDE examples.

All the PDEs in my examples “fit this setting”, for example

$$-i\partial_t u + \Delta u = |u|^2 u, \quad , \quad x \in \mathbb{T}^n$$

recalling

$$V \equiv \cup_{p \in \mathbb{R}} H_p := \left\{ u(x) = \sum_{j \in \mathbb{Z}^n} u_j e^{ij \cdot x}, \quad \sum_j \langle j \rangle^{2p} |u_j|^2 < \infty \right\}$$

$$(\langle j \rangle := \max(1, |j|))$$

we get

$$\dot{u}_j = i|j|^2 u_j + i \sum_{j_1 - j_2 + j_3 = j} u_{j_1} \bar{u}_{j_2} u_{j_3}$$

If we consider  $u_t - u_{xxt} - u_x + u_{xxx} + uu_{xxx} + 3uu_x - 2u_x u_{xx} = 0$

we get

$$\dot{u}_j = i\lambda_j u_j + P_j, \quad \lambda_j \in \mathbb{R}$$

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# Dynamical systems in $\infty$ dimension.

I shall study my system close to an elliptic fixed point that is

$$\dot{u} = F(u) = Lu + P(u)$$

where  $L$  is a (typically unbounded) linear operator with pure point purely imaginary spectrum

while  $P$  is a non-linear term (which has a zero of degree at least two at  $u = 0$ ).

Thus one can reduce to a system of the form

$$\dot{u}_j = i\lambda_j u_j + P_j(u), \quad u = \{u_j\}_{j \in I}$$

where  $\lambda_j \in \mathbb{R}$ . This is a chain of **harmonic oscillators** coupled by a **non-linearity**.

# Invariant subspaces for $\dot{u} = F(u)$

We look for a submanifold  $\mathcal{M}$  of the phase space which is **invariant** for the system and on which the **dynamics simplifies**.

Recall that  $\mathcal{M}$  is an invariant manifold for  $\dot{u} = F(u)$ , if for any initial datum  $u_0 \in \mathcal{M}$  the solution  $\Phi_t(u_0)$  stays in  $\mathcal{M}$  for all times.

$\mathcal{M}$  is  $d$  dimensional **invariant torus** if there exists a diffeomorphism

$$U : \mathbb{T}^d \rightarrow \mathcal{M} \subset \text{the phase space}, \quad \theta \mapsto U(\theta)$$

where  $\mathcal{M} = U(\mathbb{T}^d)$  is invariant...

# Invariant tori

An invariant torus

$$U : \mathbb{T}^d \rightarrow \mathcal{M} \subset \text{the phase space}, \quad \theta \mapsto U(\theta)$$

is the support of a **quasi-periodic solution of frequency  $\omega \in \mathbb{R}^d$**  if:

- The  $\omega_i$  are rationally independent
- for all  $\theta \in \mathbb{T}^d$ ,  $u(t) = U(\theta + \omega t)$  solves  $\dot{u} = F(u)$ .

# Infinite tori

One may also look for **infinite dimensional** invariant tori.

$$U : \mathbb{T}^{\mathbb{Z}} \rightarrow \mathcal{M} \subset \text{the phase space}$$

on which the dynamics is conjugated to the linear one  $\varphi \rightarrow \varphi + \omega t$  with  $\omega \in \mathbb{R}^{\mathbb{Z}}$

**We are using  $\mathbb{Z}$  as an index set because it is convenient in our examples!**

Of course for infinite tori one needs to specify the topology on  $\mathbb{T}^{\mathbb{Z}}$ . We endow  $\mathbb{T}^{\mathbb{Z}}$  with a Banach manifold structure by setting

$$\text{dist}(\theta, \varphi) = \sup_{j \in \mathbb{Z}} |\theta_j - \varphi_j|_{2\pi}$$

The topology of the phase space plays an important role!

# Perturbation Theory

We shall mostly work in a perturbative setting...

- given an **approximately invariant torus**  $\theta \rightarrow u_0(\theta)$ , are we able to prove the existence of an invariant one close by?
- assume that  $F = N_0 + P$  where  **$N_0$  admits invariant tori and  $P$  is small**...do the invariant tori persist?

Once we have found an invariant torus we would like to understand the **evolution of initial data close to the torus**.

Some interesting questions

- (Stability) Prove that all initial data which start close to the torus stay close for **finite but long times**
- (Instability) Suppose that we have two invariant tori, construct solutions which drift from one torus to the other ...estimate the time over which such drift occurs

# Small solutions

In the model

$$\dot{u}_j = i\lambda_j u_j + P_j(u), \quad u = \{u_j\}_{j \in I}$$

the simplest invariant object is  $u = 0$ .

There is a vast literature on stability/instability of elliptic fixed points!

Stability: Bourgain, Bambusi, Delort, Faou, Grebert, Szeftel, Yuan-Zhang, Cong-Mi-Wang, many more ....

Instability: Kuksin, Colliander-Keel-Staffilani-Takaoka-Tao, many more...

On PDEs without spatial confinement again a vast literature on stability/instability of zero. **very different phenomena**

# Linear theory

If take  $u$  small,  $P(u)$  is perturbative w.r.t. the linear terms...

$$\dot{u}_j = i\lambda_j u_j + P_j(u)$$

The unperturbed system is a chain of **uncoupled harmonic oscillators** with frequencies  $\lambda_j$ .

$$\dot{u}_j = i\lambda_j u_j \Rightarrow u_j(t) = e^{i\lambda_j t} u_j(0)$$

All the solutions are on **invariant tori** with dimension depending on the support of the solution

$$\mathcal{S} := \{j \in \mathbb{Z}^n : u_j(0) \neq 0\}$$

and on whether the  $\lambda_j$  are rationally independent.

$$u(t, x) = \sum_j u_j(0) e^{i\lambda_j t + i j \cdot x} = \sum_{j \in \mathcal{S}} \sqrt{\xi_j} e^{i\lambda_j t + i j \cdot x}$$

which solutions of the linear system survive the onset of the non-linearity?

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# KAM theorems in finite dim. $\dot{u}_j = i\lambda_j u_j + P_j$

In finite-dimensional nearly integrable Hamiltonian systems:

- Under some non-degeneracy assumptions (on  $\lambda$  or on  $P$ ), the majority of initial data give rise to quasi-periodic solutions that densely fill some invariant torus and are, therefore, perpetually stable.
- The KAM theorems predict persistence of most but not all quasi-periodic orbits. Obstacles arise from *small divisors* and *resonances* which should be avoided by modulating the initial data. For example for maximal KAM tori a good condition is that the frequency should be Diophantine.
- In the complementary set to the quasi-periodic orbits one may see chaotic behavior .

# KAM in infinite dimension

- In  $\infty$ -dim the scene is so far rather obscure:  
in an integrable PDE, almost-periodic solutions are typical and lie on **maximal infinite dimensional invariant tori**.  
what is their fate after perturbation? Is it still true that the majority of initial data produces perpetually stable solutions?
- There is a wide literature for existence of **quasi-periodic solutions** mostly for semilinear PDEs  
such solutions are NOT typical and correspond to lower dimensional tori
- there are very few results on infinite dimensional tori mostly for not very natural models.
- There are examples of PDEs which exhibit diffusive orbits.

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# KAM for PDEs

The first results were on model **Hamiltonian** PDEs such as the semilinear NLS with Dirichlet boundary conditions

$$-iu_t + u_{xx} + |u|^2 u + g(x, u) = 0, \quad u(t, 0) = u(t, \pi) = 0$$

- (*Semilinear PDEs with Dirichlet b.c.* : [Kuksin, Wayne, Pöschel, Kuksin-Pöschel](#), Periodic b.c. [Chierchia-You](#), [Craig-Wayne](#) '93 (periodic solutions), [Bourgain](#) '94 (quasi periodic solutions),.
- Higher dimensional manifolds:  
*Tori*: [Bourgain](#) '98,'05, [Wang](#) '10-'15, [Berti-Bolle](#) '10- '15, [Geng-You](#), [Eliasson-Kuksin](#) '10, [Geng-You-Xu](#) '10, [Procesi-P.](#) '11-'15,  
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# Some more literature: unbounded non linearities

- **semi-linear Pde's**

Kuksin '98, Kappeler-Pöeschel '03 KdV, Liu-Yuan '10,  
Zhang-Gao-Yuan '11 Hamiltonian and Reversible DNLS Berti,  
Biasco, P. , Hamiltonian and Reversible DNLW

- **Quasi-linear or fully non-linear Pde's**

**periodic solutions:** Ioss-Plotnikov-Toland '01- '10,

Alazard,Baldi capillary water waves

**quasi-periodic solutions:** Baldi, Berti, Montalto '11 Airy,Baldi,  
Berti, Haus, Montalto, Feola-Giuliani Berti-Kappeler-Montalto

**Higher Dimension:** Corsi-Montalto, Baldi-Montalto,  
Feola-Grebert

# Main difficulties. $\dot{u}_j = i\lambda_j u_j + \varepsilon P_j(u)$

Most proofs use some quadratic iterative method to construct solitons.

I will focus on three types of difficulties:

- **Resonances.** In most natural PDEs the  $\lambda_j$  are rationally dependent and thus **non generic**

The idea is to find a family of **appropriate approximately invariant tori**, typically of the form

$$\theta \rightarrow u_0(\theta) = \sum_{h=1}^d \sqrt{\xi_i} e^{i\theta_h} e^{ij_h x}$$

Where the  $\{j_1, \dots, j_d\} \subset \mathbb{Z}^n$  are **Fourier modes** and  $\xi_i$  are parameters which are modulated to deal with small divisors.

- $\{j_1, \dots, j_d\} \subset \mathbb{Z}^n$  **generic** (most choices in a big enough ball)
- $\xi_i$  are in some complicated set of positive measure in  $\mathbb{R}^d$
- **Multiplicity.**

- **Derivatives in the non-linearity** The non-linear terms can

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I will focus on three types of difficulties:

- **Resonances.**
- **Multiplicity.** In most PDEs on compact manifolds the eigenvalues  $\lambda_j$  are **multiple**...
- **Derivatives in the non-linearity** The non-linear terms can contain as many space derivatives as the linear ones.



# Main difficulties. $\dot{u}_j = i\lambda_j u_j + \varepsilon P_j(u)$

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I will focus on three types of difficulties:

- Resonances.
- Multiplicity.
- Derivatives in the non-linearity The non-linear terms can contain as many space derivatives as the linear ones.

# A result on the reversible autonomous NLS

Consider a reversible NLS equation

$$-iu_t + u_{xx} + f(u, u_x, u_{xx}) = 0 \quad (1)$$

where

$$\begin{aligned} f(u, u_x, u_{xx}) = & a_1 |u|^2 u + a_2 |u|^2 u_{xx} + a_3 |u_x|^2 u + \\ & a_4 |u_x|^2 u_{xx} + a_5 |u_{xx}|^2 u + a_6 |u_{xx}|^2 u_{xx} \end{aligned}$$

with  $a_i \in \mathbb{R}$  for  $i = 1, \dots, 6$ . Suppose that

$$(a_1, a_2, a_3, a_4, a_5, a_6) \neq (0, a, a, b, b, 0)$$

Theorem (Corsi, Feola, P. 17)

*For any generic choice of tangential sites  $j_1, \dots, j_d \in \mathbb{Z}$  and for all  $\varepsilon \in (0, \varepsilon_0)$ , there exists a positive measure set*

$$\mathcal{C}_\varepsilon \subset \varepsilon [1, 2]^d, \quad |\mathcal{C}_\varepsilon|/\varepsilon^d \rightarrow 1 \quad \text{as } \varepsilon \rightarrow 0, \quad (2)$$

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*such that for all  $\xi \in \mathcal{C}_\varepsilon$  the NLS has a quasi-periodic solution*

$$\sum_{i=1}^d \sqrt{\xi_i} e^{i\omega_i t + i j_i x} + O(\xi), \quad \omega_i(\xi) = j_i^2 + O(\xi)$$

*The quasi-periodic solutions are **analytic** and **linearly stable**.*

We also have a higher dimensional analogue.

$$-iu_t + \Delta u - F(|u|^2)u = 0$$

### Theorem (Procesi-M.P. 15)

*For any generic choice of tangential sites  $j_1, \dots, j_d \in \mathbb{Z}^n$  and for all  $\varepsilon \in (0, \varepsilon_0)$ , there exists a positive measure set*

$$\mathcal{C}_\varepsilon \subset \varepsilon [1, 2]^d, \quad |\mathcal{C}_\varepsilon|/\varepsilon^d \rightarrow \frac{1}{2} \quad \text{as } \varepsilon \rightarrow 0, \quad (3)$$

*such that for all  $\xi \in \mathcal{C}_\varepsilon$  the NLS has a quasi-periodic solution*

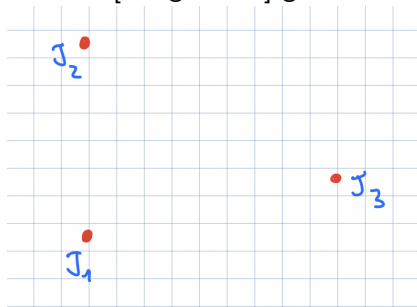
$$\sum_{i=1}^d \sqrt{\xi_i} e^{it\omega_i(\xi) + ij_i \cdot x} + O(\xi), \quad \omega_i = \omega_i(\xi) = |j_i|^2 + O(\xi)$$

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# Generic tangential sites.

In these results, the sites  $j_1, \dots, j_d \in \mathbb{Z}^n$  are generic in the sense that they **do not satisfy** a certain polynomial.

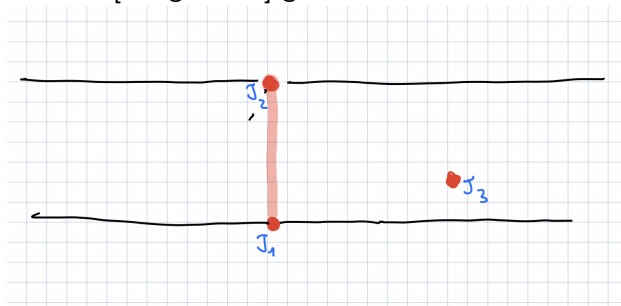
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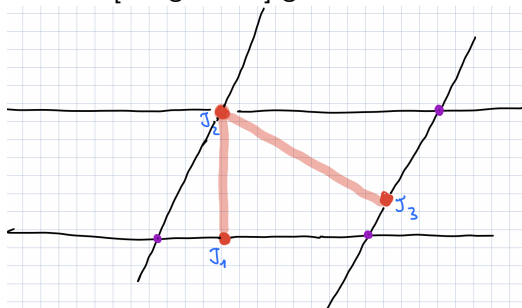
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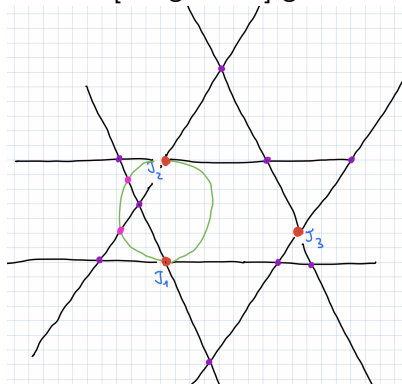
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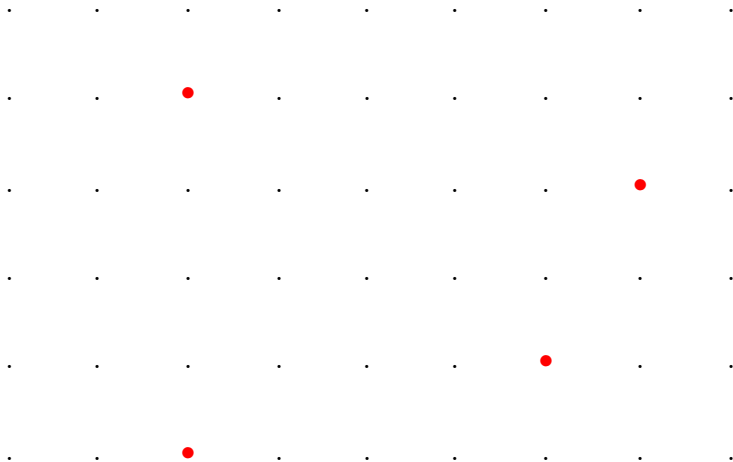
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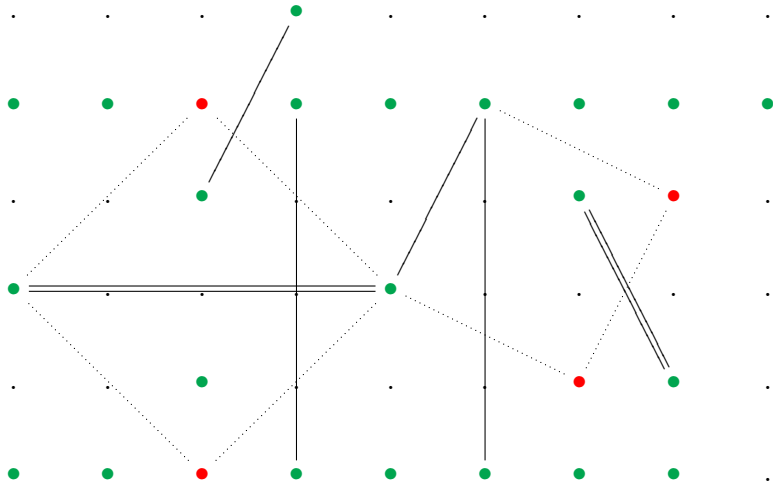




EXAMPLE:  $S$  is given by 4 points marked ●



## EXAMPLE: points connected by edges



# The Problem

## The problem

The problem consists in the study of the connected components of this graph  $\Gamma_S$  for  $S$  *generic*.

It is not hard to see that, for generic values of  $S$ , the set  $S$  is itself a connected component which we call the *special component*.

# The main combinatorial Theorem

## Theorem

*For generic choices of  $S$  the connected components of the graph  $\Gamma_S$ , different from the special component, are formed by*

*affinely independent points.*

*In particular each component has at most  $n + 1$  points.*

The proof is quite complex

it requires some algebraic geometry and a long combinatorial analysis.

# General strategies:

$$-iu_t + u_{xx} + f(u, u_x, u_{xx}) = 0, \quad -iu_t + \Delta u - F(|u|^2)u = 0,$$

- *generic* choice of tangential sites in order to avoid resonances due to  $\lambda_j \in \mathbb{Z}$ . Mostly algebraic/combinatoric methods (developed in [PP12]).
- Pseudo-differential calculus to deal with derivatives in the non-linearity (first developed for the forced Airy [BaldiBertiMontalto11])
- momentum conservation to deal with high multiplicity +Töplitz-Lipschitz property(developed in [EliassonKuksin] , [PXu]).

There is still a lot of active research on quasi-periodic solutions. Dealing with more general manifolds, quasi-linear PDEs on  $\mathbb{T}^d$ , studying stability and instability issues. 35

# Drawbacks

- Quasi-periodic solutions are **NOT typical** even for integrable equations.
- The quasi-periodic solutions described are all at least  $C^\infty$  (both in time and space)
- In order to find finite regularity solutions one needs a non-linearity with finite regularity, and even in this case the regularity is very high.

To overcome these difficulties it is natural to look at **almost-periodic solutions**.

$$U : \mathbb{T}^{\mathbb{Z}} \rightarrow \mathcal{M} \subset \text{the phase space}$$

Very few results, most on 1d NLS or NLW **with external parameters**.

$$iu_t - u_{xx} + V \star u + |u|^4 u = 0, \quad V \star u = \sum_{j \in \mathbb{Z}} V_j u_j e^{ijx}$$

$$iu_t - u_{xx} + V \star u + |u|^4 u = 0$$

Even in this special setting one is only able to construct **few** almost-periodic solutions.

**few**  $\rightsquigarrow$  **special and/or very high regularity!**

Authors	Decay of solution	Regularity of $V$
[Bo'96]	at least superexponential	analytic
[Pö'02]	at least superexponential	$\ell_2$
[Geng-Xu'12-'16]	exponential	$\ell_2$
[Bo'04 + below*]	subexponential	$\ell_\infty$
[BMP'22]	polynomial	$\ell_\infty$

below\* = [Cong-Liu-Shi-Yuan](#), [Biasco-Massetti-P.](#),  
[Cong-Mi-Shi-Wu](#), [Cong-Yuan](#) (NLW), [Cong](#), [Cong-Wu](#) high  
dimension

# Finite regularity solutions for NLS

$$iu_t + u_{xx} - V * u + F(|u|^2)u = 0 \quad (4)$$

## Theorem (Biasco-Masseti-P. 20)

For any  $p > 1$  and for *most choices of*  $V \in \ell^\infty$  there exist infinitely many (both *weak* and *classical*) almost periodic solutions

$$u(t, x) = \sum_j \hat{u}_j(t) e^{ijx}, \quad \omega_j \sim j^2, \quad \sup_j |\hat{u}_j| \langle j \rangle^p \ll 1.$$

Such solutions are **approximately supported** on **sparse subsets of  $\mathbb{Z}$** .

For example Set  $\mathcal{S} := \{j \in \mathbb{Z} : j = 2^h, \quad h \in \mathbb{N}\}$



we have a solution with  $|\hat{u}_j| \sim \langle j \rangle^{-p}$  for all  $j \in \mathcal{S}$



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we have a solution with  $|\hat{u}_j| \sim \langle j \rangle^{-p}$  *for all*  $j \in \mathcal{S}$

# Open problems

- Find Maximal tori (no conditions on  $\mathcal{S}$ ) of finite regularity.
- NLS with multiplicative potential

$$iu_t - u_{xx} + V(x)u + |u|^4 u = 0$$

- Degenerate KAM theory

$$u_{tt} + u_{xxxx} + mu + u^3 = 0$$

- Fix some (possibly very strong) regularity class, and prove that for most potentials  $V \in \ell_\infty$  typical solutions in that class are almost-periodic.

**Thanks for the attention!**